

# Classification of higher-dimensional operators in the Standard Model

Mateusz Iskrzyński

University of Warsaw

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# Higher-dimensional operators in the Standard Model

## 1 Introduction

- Effective theories
- Structure of the Standard Model

## 2 Reasoning scheme

## 3 Basis of invariant effective operators

## 4 Comparison with "Effective lagrangian analysis of new interactions and flavour conservation" by Buchmüller, Wyler (1986)

*Standard Model*  $\rightarrow$  *Extension*

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Appelquist-Carazzone decoupling theorem:

$$\mathcal{L} = \mathcal{L}_{SM}^{(4)} + \frac{1}{\Lambda} \sum_i c_i^{(5)} \mathcal{O}_i^{(5)} + \frac{1}{\Lambda^2} \sum_i c_i^{(6)} \mathcal{O}_i^{(6)} + O\left(\frac{1}{\Lambda^3}\right)$$

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22 (of 81) operators are redundant and 1 is absent.

# SM - gauge group representations structure

Field	representation (dimension)		hypercharge
	SU(3)	SU(2)	U(1)
$G_\mu$	8	1	0
$W_\mu$	1	3	0
$B_\mu$	1	1	0
$q$	3	2	$\frac{1}{6}$
$u$	3	1	$\frac{2}{3}$
$d$	3	1	$-\frac{1}{3}$
$l$	1	2	$-\frac{1}{2}$
$e$	1	1	-1
$\varphi$	1	2	$\frac{1}{2}$

$$\begin{aligned}\mathcal{L}_0 = & -\frac{1}{4}G_{\mu\nu}^A G^{A\mu\nu} - \frac{1}{4}W_{\mu\nu}^I W^{I\mu\nu} - \frac{1}{4}B_{\mu\nu} B^{\mu\nu} \\ & + (D_\mu\varphi)^\dagger(D^\mu\varphi) + m^2\varphi^\dagger\varphi - \frac{1}{2}\lambda(\varphi^\dagger\varphi)^2 \\ & + i\bar{l}\not{D}l + i\bar{e}\not{D}e + i\bar{q}\not{D}q + i\bar{u}\not{D}u + i\bar{d}\not{D}d + \\ & - (\bar{l}\Gamma_e e\varphi + \bar{q}\Gamma_u u(\tilde{\varphi}) + \bar{q}\Gamma_d d\varphi + h.c.)\end{aligned}$$

# Mass-dimension of fundamental objects in units $\hbar = c = 1$

Type	vector $V_\mu$	tensor $X_{\mu\nu}$	spinor $\Psi$	skalar $\varphi$
Dimension	$(\text{GeV})^1$	$(\text{GeV})^2$	$(\text{GeV})^{\frac{3}{2}}$	$(\text{GeV})^1$
Object	$D_\mu$	$W_{\mu\nu}, G_{\mu\nu}, B_{\mu\nu}$	$q, l, u, d, e$	$\varphi$

- ▶ for SU(3)

$$G_{\mu\nu}^A = \partial_\mu G_\nu^B - \partial_\nu G_\mu^B - g_s f^{ABC} G_\mu^B G_\nu^C$$

- ▶ for SU(2)

$$W_{\mu\nu}^I = \partial_\mu W_\nu^J - \partial_\nu W_\mu^J - g \varepsilon^{IJK} W_\mu^J W_\nu^K$$

- ▶ for U(1)

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$$

1. Description in terms of matter fields  $\varphi, \psi$ , field strength tensors  $X_{\mu\nu}$  and covariant derivatives  $D_\mu$ . Dimensional analysis.
2. Gauge and Lorentz symmetry.
3. Reduction of the set of operators using algebraic properties and SM EOM:

1. Description in terms of matter fields  $\varphi$ ,  $\psi$ , field strength tensors  $X_{\mu\nu}$  and covariant derivatives  $D_\mu$ . Dimensional analysis.  
e.g. dim-6 expressions containing both fermionic and bosonic fields:  $\psi\psi X D$ ,  $\psi\psi X\varphi$ ,  $\psi\psi\varphi\varphi\varphi$ ,  $\psi\psi\varphi\varphi D$ ,  $\psi\psi\varphi DD$ ,  $\psi\psi DDD$
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e.g.  $\psi\psi\varphi DD$ :

many possible choices of  $\psi$  - the only singlet in  $\hat{2}_{SU(2)} \otimes \hat{2}_{SU(2)}$   
hypercharge conservation

$$(q^\dagger \varepsilon \varphi^*)u, (q^\dagger \varphi)d, (l^\dagger \varphi)e, + h.c.$$

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# Reasoning scheme

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Lorentz structure contains 2 singlets:

$$\begin{aligned} \left(\frac{1}{2}, 0\right) \otimes \left(\frac{1}{2}, 0\right) \otimes \left(\frac{1}{2}, \frac{1}{2}\right) \otimes \left(\frac{1}{2}, \frac{1}{2}\right) &= (0, 0) \oplus (0, 0) \oplus (1, 0) \oplus (2, 0) \\ &\oplus (1, 1) \oplus (0, 1) \oplus (1, 1) \oplus (2, 1) \end{aligned}$$

3. Reduction of the set of operators using algebraic properties and SM EOM:

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2 independent Lorentz invariants (for each):

$$\bar{\psi}_L \psi_R \varphi D_\mu D^\mu \quad \bar{\psi}_L \sigma_{\mu\nu} \psi_R \varphi D^\mu D^\nu$$

3. Reduction of the set of operators using algebraic properties and SM EOM:

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We have (omitting full div) the following operators:

$$(\bar{\psi}_L \sigma_{\mu\nu} \psi_R)(D^\mu D^\nu \varphi) \quad (1)$$

$$(\bar{\psi}_L \sigma_{\mu\nu} D^\mu D^\nu \psi_R) \varphi \quad (2)$$

$$(\bar{\psi}_L \sigma_{\mu\nu} D^\mu \psi_R)(D^\nu \varphi) \quad (3)$$

$$(\bar{\psi}_L D_\mu D^\mu \psi_R) \varphi \quad (4)$$

$$(\bar{\psi}_L D_\mu \psi_R)(D^\mu \varphi) \quad (5)$$

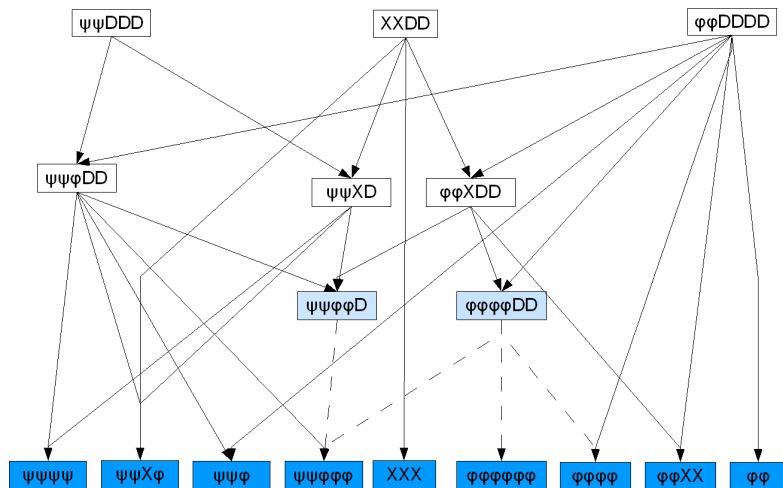
$$(\bar{\psi}_L \psi_R)(D_\mu D^\mu \varphi) \quad (6)$$

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We can reduce:

$$\begin{aligned}(\bar{\psi}_L \sigma_{\mu\nu} \psi_R)(D^\mu D^\nu \varphi) &= \frac{1}{2}(\bar{\psi}_L \sigma_{\mu\nu} \psi_R)([D^\mu, D^\nu] \varphi) \\ &= \frac{1}{2}(\bar{\psi}_L \sigma_{\mu\nu} \psi_R)(ig\mathcal{W}^{\mu\nu} + ig'\mathcal{B}^{\mu\nu})\varphi \sim \boxed{\psi\psi X\varphi}\end{aligned}$$

# Reduction scheme



# Bosonic invariant operators

$\varphi^6$	XXX	$\varphi\varphi XX$	$\varphi^4 DD$
$(\varphi^\dagger\varphi)^3$	$\epsilon^{IJK} W_\mu^I{}^\nu W_\nu^J{}^\rho W_\rho^K{}^\mu$ $\epsilon^{IJK} \tilde{W}^{I\mu\nu} W_{\nu\delta}^J W_\mu^{K\delta}$ $f^{ABC} G_\mu^A{}^\nu G_\nu^B{}^\rho G_\rho^C{}^\mu$ $f^{ABC} \tilde{G}_\mu^A{}^\nu G_\nu^B{}^\delta G_\delta^C{}^\mu$	$\varphi^\dagger T^I \varphi W_{\mu\nu}^I B^{\mu\nu}$ $\varphi^\dagger T^I \varphi W_{\mu\nu}^I \tilde{B}^{\mu\nu}$ $\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$ $\varphi^\dagger \varphi W_{\mu\nu}^I \tilde{W}^{I\mu\nu}$ $\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$ $\varphi^\dagger \varphi G_{\mu\nu}^A \tilde{G}^{A\mu\nu}$ $\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$ $\varphi^\dagger \varphi B_{\mu\nu} \tilde{B}^{\mu\nu}$	$(\varphi^\dagger\varphi)(D_\mu\varphi)^\dagger(D^\mu\varphi)$ $[\varphi^\dagger(D^\mu\varphi)][(D_\mu\varphi)^\dagger\varphi]$

# Invariant operators with 2 fermions

$\psi\psi\varphi\varphi D$	$\psi\psi\varphi\varphi\varphi$	$\psi\psi X\varphi$
$(\bar{q}\gamma^\mu q)(\varphi^\dagger D_\mu\varphi)$	$[\bar{u}(\tilde{\varphi}^\dagger q)](\varphi^\dagger\varphi)$	$\bar{d}\sigma^{\mu\nu}\lambda^A(\varphi^\dagger q)G_{\mu\nu}^A$
$(\bar{q}\gamma^\mu T^I q)(\varphi^\dagger T^I D_\mu\varphi)$	$[\bar{d}(\varphi^\dagger q)](\varphi^\dagger\varphi)$	$\bar{u}\sigma^{\mu\nu}\lambda^A(\tilde{\varphi}^\dagger q)G_{\mu\nu}^A$
$(\bar{u}\gamma^\mu d)(\varphi^\dagger D_\mu\tilde{\varphi})$	$[\bar{e}(\varphi^\dagger l)](\varphi^\dagger\varphi)$	$\bar{d}\sigma^{\mu\nu}T^I(\varphi^\dagger q)W_{\mu\nu}^I$
$(\bar{u}\gamma^\mu u)(\varphi^\dagger D_\mu\varphi)$		$\bar{u}\sigma^{\mu\nu}T^I(\tilde{\varphi}^\dagger q)W_{\mu\nu}^I$
$(\bar{d}\gamma^\mu d)(\varphi^\dagger D_\mu\varphi)$		$\bar{e}\sigma^{\mu\nu}T^I(\varphi^\dagger l)W_{\mu\nu}^I$
$(\bar{e}\gamma^\mu e)(\varphi^\dagger D_\mu\varphi)$		$\bar{u}\sigma^{\mu\nu}(\tilde{\varphi}^\dagger q)B_{\mu\nu}$
$(\bar{l}\gamma^\mu l)(\varphi^\dagger D_\mu\varphi)$		$\bar{d}\sigma^{\mu\nu}(\varphi^\dagger q)B_{\mu\nu}$
$(\varphi^\dagger l)\gamma^\mu(\bar{l}D_\mu\varphi)$		$\bar{e}\sigma^{\mu\nu}(\varphi^\dagger l)B_{\mu\nu}$

# Fermionic operators

$\bar{L}\bar{L}\bar{L}\bar{L}$	$\bar{R}\bar{R}\bar{R}\bar{R}$
$(\bar{l}_{p_1} \gamma_\mu l_{p_2})(\bar{l}_{p_3} \gamma^\mu l_{p_4})$	$(\bar{e} \gamma_\mu e)(\bar{e} \gamma^\mu e)$
$(\bar{q}_{p_1} \gamma_\mu q_{p_2})(\bar{q}_{p_3} \gamma^\mu q_{p_4})$	$(\bar{u}_{p_1} \gamma_\mu u_{p_2})(\bar{u}_{p_3} \gamma^\mu u_{p_4})$
$(\bar{q}_{p_1} \gamma_\mu T^I q_{p_2})(\bar{q}_{p_3} \gamma^\mu T^I q_{p_4})$	$(\bar{d}_{p_1} \gamma_\mu d_{p_2})(\bar{d}_{p_3} \gamma^\mu d_{p_4})$
$(\bar{q}_{p_1} \gamma_\mu q_{p_2})(\bar{l}_{p_3} \gamma^\mu l_{p_4})$	$(\bar{u}_{p_1} \gamma_\mu u_{p_2})(\bar{e} \gamma^\mu e)$
$(\bar{q}_{p_1} \gamma_\mu T^I q_{p_2})(\bar{l}_{p_3} \gamma^\mu T^I l_{p_4})$	$(\bar{d}_{p_1} \gamma_\mu d_{p_2})(\bar{e} \gamma^\mu e)$
	$(\bar{u}_{p_1} \gamma_\mu u_{p_2})(\bar{d}_{p_3} \gamma^\mu d_{p_4})$
	$(\bar{u}_{p_1} \gamma_\mu T^A u_{p_2})(\bar{d}_{p_3} T^A d_{p_4})$



# Fermionic operators

$\bar{L}R\bar{R}L$	$\bar{L}R\bar{L}R$
$(\bar{l}e)(\bar{e}l)$	$(\bar{q}_{p_1} u_{p_2})\varepsilon(\bar{q}_{p_3} d_{p_4})$
$(\bar{u}l)(\bar{l}u)$	$(\bar{q}_{p_1} T^A u_{p_2})\varepsilon(\bar{q}_{p_3} T^A d_{p_4})$
$(\bar{d}l)(\bar{l}d)$	$(\bar{q}u)\varepsilon(\bar{l}e)^T$
$(\bar{q}_{p_1} e_{p_2})(\bar{e}_{p_3} q_{p_4})$	$(\bar{q}e)\varepsilon(\bar{l}u)^T$
$(\bar{q}_{p_1} u_{p_2})(\bar{u}_{p_3} q_{p_4})$	
$(\bar{q}_{p_1} T^A u_{p_2})(\bar{u}_{p_3} T^A q_{p_4})$	
$(\bar{q}_{p_1} d_{p_2})(\bar{d}_{p_3} q_{p_4})$	
$(\bar{q}_{p_1} T^A d_{p_4})(\bar{d}_{p_3} T^A q_{p_2})$	
$(\bar{q}d)(\bar{e}l)$	

$$\frac{1}{2}\partial_\mu(\varphi^\dagger\varphi)\partial^\mu(\varphi^\dagger\varphi)$$

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$$(\bar{\psi}_R D_\mu \psi_L)(D^\mu \varphi)$$

$$[(D_\mu \bar{\psi}_R)\psi_L](D^\mu \varphi)$$

$$(\bar{\psi}\gamma_\mu D_\nu \psi)X^{\mu\nu}$$

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$$(\bar{l}_{p1}\gamma_\mu T^I l_{p2})(\bar{l}_{p3}\gamma^\mu T^I l_{p4})$$

$$T_{ij}^I T_{kl}^I = \frac{1}{2}\delta_{il}\delta_{kj} - \frac{1}{2N}\delta_{ij}\delta_{kl}$$

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$$T^I_{ij} T^I_{kl} = \frac{1}{2}\delta_{il}\delta_{kj} - \frac{1}{2N}\delta_{ij}\delta_{kl}$$

The absent one:  $(\bar{q}e)\varepsilon(\bar{l}u)^T$

$$(l^T \tilde{\varphi}^*)(\tilde{\varphi}^\dagger l) \quad (1)$$

$$\varepsilon^{abc} (l_{p_1}^T \gamma_\mu u^b)_{p_2} \varepsilon (d_{p_3}^{cT} \gamma^\mu q_{p_4}^a) \quad (2)$$

$$\varepsilon^{abc} (e_{p_1}^T \gamma_\mu q_{p_2}^{ai}) \varepsilon^{ij} (u_{p_3}^{cT} \gamma^\mu q_{p_4}^{bj}) \quad (3)$$

$$\varepsilon^{abc} (e_{p_1}^T u_{p_4}^a) (u_{p_3}^{bT} d_{p_2}^c) \quad (4)$$

$$\varepsilon^{abc} (l_{p_1}^T \varepsilon q_{p_2}^a) (q_{p_3}^{bT} \varepsilon q_{p_4}^c) \quad (5)$$

$$\varepsilon^{abc} (l_{p_1}^T \varepsilon \tau^I q_{p_2}^a) (q_{p_3}^{bT} \varepsilon \tau^I q_{p_4}^c) \quad (6)$$

# The importance of classification

How to discover new physics?

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Recent papers with redundant operators:

- ▶ J. A. Aguilar-Saavedra, "Single top quark production at LHC with anomalous  $Wtb$  couplings", Nucl. Phys. B804 (2008) 160;
- ▶ K. Agashe, R. Contino, "Composite Higgs-mediated flavor-changing neutral current", Phys. Rev. D 80, 075016 (2009);
- ▶ S. Kanemura, K. Tsumura, "Effects of the anomalous Higgs couplings on the Higgs boson production at the Large Hadron Collider", Eur. Phys. J. C63 (2009) 11;



# Questions



# Example: reduction of $(\bar{\psi}D_\mu\psi)(D^\mu\varphi)$

$$\begin{aligned}
 (\bar{\psi}D_\mu\psi)(D^\mu\varphi) &= (\bar{\psi}D_\nu\eta^{\nu\mu}\psi)(D_\mu\varphi) = (\bar{\psi}D_\nu\frac{1}{2}\{\gamma^\nu, \gamma^\mu\}\psi)(D_\mu\varphi) \\
 &= \frac{1}{2}(\bar{\psi}D_\nu\gamma^\nu\gamma^\mu\psi)(D_\mu\varphi) + \frac{1}{2}(\bar{\psi}\gamma^\mu\underline{D}\psi)(D_\mu\varphi) \\
 &= \boxed{\psi\psi\varphi\varphi D} + \frac{1}{2}D_\nu[(\bar{\psi}\gamma^\nu\gamma^\mu\psi)(D_\mu\varphi)] - \frac{1}{2}(\bar{\psi}\overleftarrow{D}\gamma^\mu\psi)(D_\mu\varphi) + \\
 &\quad - \frac{1}{2}(\bar{\psi}\gamma^\nu\gamma^\mu\psi)(D_\nu D_\mu\varphi)
 \end{aligned}$$

but

$$(\bar{\psi}\gamma^\nu\gamma^\mu\psi)(D_\nu D_\mu\varphi) = (\bar{\psi}(-i)\sigma^{\nu\mu}\psi)(\sum_k X_{\nu\mu}^k\varphi) + (\bar{\psi}\psi)(\underline{D}^\mu D_\mu\varphi)$$

so

$$(\bar{\psi}_L D_\mu \psi_R)(D^\mu \varphi) = \boxed{\psi\psi\varphi\varphi D} + \boxed{\psi\psi X\varphi} + \boxed{\psi\psi\varphi} + \boxed{\psi\psi\varphi\varphi} + \boxed{\psi\psi\psi\psi}$$

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 (\bar{\psi} D_\mu \psi)(D^\mu \varphi) &= (\bar{\psi} D_\nu \eta^{\nu\mu} \psi)(D_\mu \varphi) = (\bar{\psi} D_\nu \frac{1}{2} \{\gamma^\nu, \gamma^\mu\} \psi)(D_\mu \varphi) \\
 &= \frac{1}{2} (\bar{\psi} D_\nu \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi) + \frac{1}{2} (\bar{\psi} \gamma^\mu \overleftarrow{D} \psi)(D_\mu \varphi) \\
 &= \boxed{\psi \psi \varphi \varphi D} + \frac{1}{2} D_\nu [(\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi)] - \frac{1}{2} (\bar{\psi} \overleftarrow{D} \gamma^\mu \psi)(D_\mu \varphi) + \\
 &\quad - \frac{1}{2} (\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\nu D_\mu \varphi)
 \end{aligned}$$

but

$$(\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\nu D_\mu \varphi) = (\bar{\psi} (-i) \sigma^{\nu\mu} \psi) \left( \sum_k X_{\nu\mu}^k \varphi \right) + (\bar{\psi} \psi) (\underline{D^\mu D_\mu \varphi})$$

so

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 &= \frac{1}{2} (\bar{\psi} D_\nu \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi) + \frac{1}{2} (\bar{\psi} \gamma^\mu \overleftarrow{D} \psi)(D_\mu \varphi) \\
 &= \boxed{\psi \psi \varphi \varphi D} + \frac{1}{2} D_\nu [(\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi)] - \frac{1}{2} (\bar{\psi} \overleftarrow{D} \gamma^\mu \psi)(D_\mu \varphi) + \\
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$$= \frac{1}{2} (\bar{\psi} D_\nu \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi) + \frac{1}{2} (\bar{\psi} \gamma^\mu \overleftarrow{D} \psi)(D_\mu \varphi)$$

$$= \boxed{\psi \psi \varphi \varphi D} + \frac{1}{2} D_\nu [(\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi)] - \frac{1}{2} (\bar{\psi} \overleftarrow{D} \gamma^\mu \psi)(D_\mu \varphi) + \\ - \frac{1}{2} (\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\nu D_\mu \varphi)$$

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# Example: reduction of $(\bar{\psi} D_\mu \psi)(D^\mu \varphi)$

$$\begin{aligned}
 (\bar{\psi} D_\mu \psi)(D^\mu \varphi) &= (\bar{\psi} D_\nu \eta^{\nu\mu} \psi)(D_\mu \varphi) = (\bar{\psi} D_\nu \frac{1}{2} \{\gamma^\nu, \gamma^\mu\} \psi)(D_\mu \varphi) \\
 &= \frac{1}{2} (\bar{\psi} D_\nu \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi) + \frac{1}{2} (\bar{\psi} \gamma^\mu \overleftarrow{D} \psi)(D_\mu \varphi) \\
 &= \boxed{\psi \psi \varphi \varphi D} + \frac{1}{2} D_\nu [(\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\mu \varphi)] - \frac{1}{2} (\bar{\psi} \overleftarrow{D} \gamma^\mu \psi)(D_\mu \varphi) + \\
 &\quad - \frac{1}{2} (\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\nu D_\mu \varphi)
 \end{aligned}$$

but

$$(\bar{\psi} \gamma^\nu \gamma^\mu \psi)(D_\nu D_\mu \varphi) = (\bar{\psi} (-i) \sigma^{\nu\mu} \psi) \left( \sum_k X_{\nu\mu}^k \varphi \right) + (\bar{\psi} \psi) (\underline{D^\mu D_\mu \varphi})$$

so

$$(\bar{\psi}_L D_\mu \psi_R)(D^\mu \varphi) = \boxed{\psi \psi \varphi \varphi D} + \boxed{\psi \psi X \varphi} + \boxed{\psi \psi \varphi} + \boxed{\psi \psi \varphi \varphi} + \boxed{\psi \psi \psi \psi}$$